A concrete example

Gibbs principle

Stochast

Reterence:

# Gibbs principle on path space and relations with stochastic control

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### A concrete example

# A concrete example

Conditioning Entropy

Gibbs principle

Stochasti control

Reference



From "Modern Times", Charlie Chaplin, 1936, 13th minute.

 $\Rightarrow$  The Tramp works.

# Let us model *The Tramp*'s work

# A concrete example

Conditionir

Gibbs principle

Stochasti control

References

### Three possible states:



The Tramp is on time.



The Tramp takes a break.



The Tramp is late.

### Random disturbance

# A concrete example

Conditio

Gibbs principle

Stochasti control

References

#### Deterministic evolution:

$$\mathsf{Tramp}_{t+1} = F(t, \mathsf{Tramp}_t).$$

 $\hookrightarrow$  It is sufficient to know Tramp<sub>t=0</sub>.

### Random disturbance

# A concrete example

Entropy

Gibbs principle

Stochasti control

Reference

#### Deterministic evolution:

$$\mathsf{Tramp}_{t+1} = F(t, \mathsf{Tramp}_t).$$

 $\hookrightarrow$  It is sufficient to know Tramp<sub>t=0</sub>.

#### But...



⇒ F does not account for interaction.

### Probabilistic model

A concrete example

Conditio

Gibbs principle

Stochastic control

References

Let us a consider a prior distribution

$$\nu \in \mathcal{P}(\{\text{on time}, \text{late}, \text{break}\}),$$

for instance

$$\nu$$
(on time) =  $\nu$ (late) =  $\nu$ (break) =  $\frac{1}{3}$ .

### Probabilistic model

A concrete example

Entropy

Gibbs principle

Stochasti control

References

Let us a consider a prior distribution

$$\nu \in \mathcal{P}(\{\mathsf{on\ time},\mathsf{late},\mathsf{break}\}),$$

for instance

$$\nu$$
(on time) =  $\nu$ (late) =  $\nu$ (break) =  $\frac{1}{3}$ .

This is not accurate...



...because *The Tramp* is never on time!

### Adding information

A concrete example

Condition

Entropy

Gibbs principle

Stochasti control

References

#### Conditioning

$$u_{\text{on time}} := \nu(\cdot|\mathsf{Tramp} \neq \mathsf{on time}) = Z^{-1} \mathbb{1}_{\{\mathsf{Tramp} \neq \mathsf{on time}\}} \nu$$

- $\rightarrow \nu_{\text{on time}}$  matches observation:  $\nu_{\text{on time}}$  (on time) = 0.
- $\hookrightarrow$  Is this choice of  $\nu_{\text{on time}}$  optimal?...

### Adding information

A concrete example

Conditio

Gibbs principle

Stochastic control

References

### Conditioning

$$\nu_{\text{on-time}} := \nu(\cdot|\mathsf{Tramp} \neq \mathsf{on time}) = Z^{-1} \mathbb{1}_{\{\mathsf{Tramp} \neq \mathsf{on time}\}} \nu$$

- $\rightarrow \nu_{\text{on time}}$  matches observation:  $\nu_{\text{on time}}$  (on time) = 0.
- $\hookrightarrow$  Is this choice of  $\nu_{\text{on time}}$  optimal?...

...Yes, it is!

$$\nu_{\text{on time}} = \underset{\substack{\mu \in \mathcal{P} \\ \mu \text{(on time)} = 0}}{\operatorname{argmin}} H(\mu|\nu),$$

where 
$$H(\mu|\nu) \ge 0$$
 and  $H(\mu|\nu) \Leftrightarrow \mu = \nu$ .

⇒ Natural distance w.r.t. conditioning.

### Relative entropy

A concrete example

Conditioning

Entropy

Gibbs principle

Stochastic control

References

Given  $\nu \in \mathcal{P}(\mathbb{R}^d)$ , the relative entropy is

$$H(\mu|\nu) := egin{cases} \int_{\mathcal{E}} \log rac{\mathrm{d}\mu}{\mathrm{d}\nu} \mathrm{d}\mu & ext{if } \mu \ll \nu, \ +\infty & ext{otherwise}. \end{cases}$$

Given  $A \subset \mathbb{R}^d$  with  $\nu(A) > 0$ ,

# Relative entropy

A concrete example

Conditioning

Entropy

Gibbs principle

Stochastic control

References

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Given  $A \subset \mathbb{R}^d$  with  $\nu(A) > 0$ ,

Proof.

$$H(\mu|\nu) + \mu(A) = H(\nu_{A^c}|\nu) + \nu(A) + \underbrace{H(\mu|\nu_{A^c})}_{> 0}.$$

# Adding statistical information

A concrete example

F .

Gibbs principle

Stochasti control

References

### Defining

$$\psi(\mathsf{Tramp}) := \mathbb{1}_{\{\mathsf{Tramp} = \mathsf{on time}\}} - \frac{1}{42},$$

we now measure that

$$\mathbb{E}\,\psi(\mathsf{Tramp}) \leq \mathsf{0}.$$

 $\hookrightarrow$  How can we correct  $\mu$  to account for this?



### Adding statistical information

A concrete

Gibbs principle

Stochasti control

References

### Defining

$$\psi(\mathsf{Tramp}) := \mathbbm{1}_{\{\mathsf{Tramp} = \mathsf{on time}\}} - \frac{1}{42},$$

we now measure that

$$\mathbb{E}\,\psi(\mathsf{Tramp}) \leq \mathsf{0}.$$

 $\hookrightarrow$  How can we correct  $\mu$  to account for this?



A natural candidate is

$$\underset{\substack{\mu \in \mathcal{P} \\ \langle \mu, \psi \rangle \leq 0}}{\operatorname{argmin}} \ H(\mu|\nu),$$

where  $\langle \mu, \psi \rangle := \int \psi \, \mathrm{d}\mu$ .

### Gibbs variational principle

A concrete example

Gibbs principle

Statistical physics

Path space

Stochasti control

References

### Theorem (Gibbs variational principle)

Given  $\nu$  in  $\mathcal{P}(\mathbb{R}^d)$  and  $\psi: \mathbb{R}^d \to \mathbb{R}$  continuous, bounded from below,

$$\inf_{\substack{\mu \in \mathcal{P}(\mathbb{R}^d) \\ \langle \mu, \psi \rangle \le 0}} H(\mu|\nu)$$

is realised by a unique measure  $\mu_{\overline{\beta}}$  for some  $\overline{\beta} \in \mathbb{R}$ , where

$$\frac{\mathrm{d}\mu_{\overline{\beta}}}{\mathrm{d}\nu}(x) = Z_{\overline{\beta}}^{-1} e^{-\overline{\beta}\psi(x)}.$$

- $\hookrightarrow$  Single linear constraint [SZ91; DZ96].
- $\hookrightarrow$  Lagrange multiplier  $\overline{\beta}$ .

# Proof of the Gibbs principle in $\mathbb{R}^d$

A concrete example

#### Gibbs principle

Statistical physics

Path space

Stochasti control

References

#### Proof.

The Gibbs free energy

$$G(\beta) := \inf_{\mu \in \mathcal{P}(\mathbb{R}^d)} H(\mu|\nu) + \beta \langle \mu, \psi \rangle,$$

is uniquely realised by  $\mu = \mu_{\beta}$ , because

$$\geq 0$$

$$H(\mu|\nu) + \beta\langle\mu,\psi\rangle = H(\mu_{\beta}|\nu) + \beta\langle\mu_{\beta},\psi\rangle + H(\mu|\mu_{\beta}),$$

hence 
$$G(\beta) = -\log Z_{\beta}$$
.

# Proof of the Gibbs principle in $\mathbb{R}^d$

A concrete example

#### Gibbs principle

Statistical physic

Path space

Stochasti control

References

#### Proof.

The Gibbs free energy

$$G(\beta) := \inf_{\mu \in \mathcal{P}(\mathbb{R}^d)} H(\mu|\nu) + \beta \langle \mu, \psi \rangle,$$

is uniquely realised by  $\mu = \mu_{\beta}$ , because

$$\geq 0$$

$$H(\mu|\nu) + \beta\langle\mu,\psi\rangle = H(\mu_{\beta}|\nu) + \beta\langle\mu_{\beta},\psi\rangle + H(\mu|\mu_{\beta}),$$

hence  $G(\beta) = -\log Z_{\beta}$ . Moreover,

$$\inf_{\substack{\mu \in \mathcal{P}(\mathbb{R}^d) \\ \langle \mu, \psi \rangle \leq 0}} H(\mu|\nu) = \inf_{\substack{\mu \in \mathcal{P}(\mathbb{R}^d) \\ \langle \mu, \psi \rangle \leq 0}} \sup_{\beta \geq 0} H(\mu|\nu) + \beta \langle \mu, \psi \rangle$$
$$\geq \sup_{\alpha \in \mathbb{R}^n} \frac{G(\beta)}{\sigma},$$

and  $\beta \mapsto \langle \mu_{\beta}, \psi \rangle$  is continuous and decreasing. . .

# An alternative approach to entropy

A concrete example

Gibbs principle

Large deviations

Path space

Stochasti

References

#### Law of large numbers

For  $(X^i)_{i\geq 1}$  independent  $\nu$ -distributed variables,

$$\frac{1}{N}\sum_{i=1}^{N}\varphi(X^{i})\xrightarrow[N\to+\infty]{\text{a.s.}}\mathbb{E}\varphi(X^{1})=\langle\nu,\varphi\rangle$$

for every  $\varphi$ , so that

$$\pi(\vec{X}^N) := \frac{1}{N} \sum_{i=1}^N \delta_{X^i} \xrightarrow[N \to +\infty]{\text{weak}} \nu.$$

# An alternative approach to entropy

A concrete example

Gibbs principle

Large deviation

Statistical physics

Path space

Stochastic control

References

#### Law of large numbers

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for every  $\varphi$ , so that

$$\pi(\vec{X}^N) := \frac{1}{N} \sum_{i=1}^N \delta_{X^i} \xrightarrow[N \to +\infty]{\text{weak}} \nu.$$

$$\mathbb{P}(\pi(\vec{X}^N) = \mu) = ?$$

# An alternative approach to entropy

A concrete example

Gibbs principle

Large deviation

Statistical physic

Path space

Stochastic control

References

#### Law of large numbers

For  $(X^i)_{i\geq 1}$  independent  $\nu$ -distributed variables,

$$\frac{1}{N}\sum_{i=1}^{N}\varphi(X^{i})\xrightarrow[N\to+\infty]{\text{a.s.}}\mathbb{E}\varphi(X^{1})=\langle\nu,\varphi\rangle$$

for every  $\varphi$ , so that

$$\pi(\vec{X}^N) := \frac{1}{N} \sum_{i=1}^N \delta_{X^i} \xrightarrow[N \to +\infty]{\text{weak}} \nu.$$

→ What about fluctuations?

$$\mathbb{P}(\pi(\vec{X}^N) \simeq \mu) \simeq e^{-NH(\mu|\nu)}$$

 $\Rightarrow$  H quantifies deviations from the LLN.

# Gibbs principle through large deviations

A concrete example

Gibbs principle

Large deviations

Statistical physic

Path space

Stochast control

References

From large deviations,

$$\mathbb{P}(\pi(\vec{X}^N) \in A) \simeq \exp\left[-N\inf_{\mu \in A} H(\mu|\nu)\right],$$

so that

$$\mathbb{P}\big(\langle \pi(\vec{X}^{N}), \psi \rangle \leq 0\big) \simeq \exp\left[-N\inf_{\mu, \langle \mu, \psi \rangle \leq 0} H(\mu|\nu)\right].$$

# Gibbs principle through large deviations

A concrete example

Gibbs principle

Large deviations

Statistical physic

Path space

Stochasti control

References

From large deviations,

$$\mathbb{P}(\pi(\vec{X}^N) \in A) \simeq \exp\left[-N\inf_{\mu \in A} H(\mu|\nu)\right],$$

so that

$$\mathbb{P}\big(\langle \pi(\vec{X}^{N}), \psi \rangle \leq 0\big) \simeq \exp\bigg[-N\inf_{\mu, \langle \mu, \psi \rangle \leq 0} H(\mu|\nu)\bigg].$$

Consequently,

$$\begin{split} \mathbb{P}\big( \left< \pi(\vec{X}^N) \in A \, | \, \left< \pi(\vec{X}^N), \psi \right> \leq 0 \, \big) \simeq \\ \exp \left[ -N \bigg( \inf_{\mu \in A, \langle \mu, \psi \rangle \leq 0} H(\mu|\nu) - \inf_{\mu, \langle \mu, \psi \rangle \leq 0} H(\mu|\nu) \bigg) \right]. \end{split}$$

### Gibbs principle in statistical mechanics

A concrete example

Gibbs principle

Statistical physics

Path space

Stochasti control

References

#### Canonical ensemble

For the kinetic energy  $\psi(\mathbf{v}) := \frac{1}{2} m \mathbf{v}^2$  in  $\mathbb{R}^3$ ,

$$\begin{split} \big(\,\pi(\vec{V}^{N}) \in A \,|\, \langle \pi(\vec{V}^{N}), \psi \rangle &= \overline{\psi}\,\big) \xrightarrow[N \to +\infty]{\mathsf{law}} \\ & \underset{\mu \in \mathcal{P}(\mathbb{R}^{3}), \, \langle \mu, \psi \rangle = \overline{\psi}}{\mathsf{argmax}} \,- \int_{\mathbb{R}^{3}} \mu(v) \log \mu(v) \mathrm{d}v, \end{split}$$

### Gibbs principle in statistical mechanics

A concrete example

Gibbs principle

Statistical phys

Path space

Stochastic control

References

#### Canonical ensemble

For the kinetic energy  $\psi(v):=rac{1}{2}mv^2$  in  $\mathbb{R}^3$ ,

$$(\pi(\vec{V}^N) \in A \, | \, \langle \pi(\vec{V}^N), \psi \rangle = \overline{\psi} \, ) \xrightarrow[N \to +\infty]{\text{law}}$$

$$\underset{\mu \in \mathcal{P}(\mathbb{R}^3), \, \langle \mu, \psi \rangle = \overline{\psi}}{\text{argmax}} - \int_{\mathbb{R}^3} \mu(v) \log \mu(v) \mathrm{d}v,$$

giving the Maxwell-Boltzmann distribution

$$\overline{\mu}(v) := \left(\frac{m}{2\pi k_B T}\right)^{3/2} \exp\left[-\frac{mv^2}{2k_B T}\right],$$

where 
$$\overline{\psi} = \frac{3}{2}k_BT = \frac{3}{2\beta}$$
.

# Infinitely many constraints

A concrete

Gibbs principle

Statistical physic

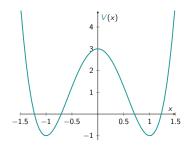
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Stochast control

References

Example of i.i.d. diffusion processes in  $\mathbb{R}^d$ 

$$dX_t^{i,N} = -\nabla V(X_t^{i,N})dt + dB_t^{i,N}, \quad 1 \le i \le N,$$



Confinement potential V

with mean-field conditioning:

$$\forall t \in [0, T], \quad \frac{1}{N} \sum_{i=1}^{N} V(X_t^{i,N}) \leq 0.$$

# General setting

A concrete example

Gibbs principle

Statistical physic

Path space

Stochasti control

References

Let  $(X_{[0,T]}^i)_{i\geq 1}$  be independent  $\nu_{[0,T]}$ -distributed  $C([0,T],\mathbb{R}^d)$ -valued variables, and let

$$\pi(\vec{X}_t^N) := \frac{1}{N} \sum_{i=1}^N \delta_{X_t^i},$$

denote the empirical measure at time t.

We condition by

$$\forall t \in [0, T], \quad \Psi(\pi(\vec{X}_t^N)) \le 0. \tag{1}$$

# General setting

A concrete example

Gibbs principle

Large deviations

Statistical physics

Path space

Stochasti control

References

Let  $(X_{[0,T]}^i)_{i\geq 1}$  be independent  $\nu_{[0,T]}$ -distributed  $C([0,T],\mathbb{R}^d)$ -valued variables, and let

$$\pi(\vec{X}_t^N) := \frac{1}{N} \sum_{i=1}^N \delta_{X_t^i},$$

denote the empirical measure at time t.

We condition by

$$\forall t \in [0, T], \quad \Psi(\pi(\vec{X}_t^N)) \le 0. \tag{1}$$

As  $N \to +\infty$ ,  $\operatorname{Law}(X^1_{[0,T]}|\ (1))$  converges towards

$$\underset{\substack{\mu_{[0,T]} \in \mathcal{P}(\mathcal{C}([0,T],\mathbb{R}^d))\\ \forall t \in [0,T], \ \Psi(\mu_t) \leq 0}}{\operatorname{argmin}} \frac{H(\mu_{[0,T]}|\nu_{[0,T]}).$$

### Regularity assumptions

A concrete example

Gibbs principle

Statistical physic

Path space

Stochasti control

References

The map  $\Psi: \mathcal{P}(\mathbb{R}^d) \to \mathbb{R}$  is  $C^1$  in the following sense: for any  $\mu, \mu' \in \mathcal{P}(\mathbb{R}^d)$ ,

$$\frac{\mathrm{d}}{\mathrm{d}\varepsilon}\bigg|_{\varepsilon=0}\Psi(\mu+\varepsilon(\mu'-\mu))=\langle\mu-\mu',\frac{\delta\Psi}{\delta\mu}(\mu)\rangle,$$

for a continuous  $\frac{\delta\Psi}{\delta\mu}:\mathcal{P}(\mathbb{R}^d) imes\mathbb{R}^d o\mathbb{R}.$ 

# Regularity assumptions

A concrete example

Gibbs principle

Large deviations

Statistical physics

Path space

Stochastic control

References

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for a continuous  $\frac{\delta\Psi}{\delta\mu}:\mathcal{P}(\mathbb{R}^d) imes\mathbb{R}^d o\mathbb{R}.$ 

### Example

The constraints

$$\Psi_{\mathcal{V}}(\mu) = \langle \mu, \mathcal{V} \rangle$$
 and  $\Psi_{\mathcal{W}}(\mu) = \langle \mu, \mathcal{W} \star \mu \rangle$ 

are respectively linear and non-convex in  $\mu$ . For even W,

$$\frac{\delta \Psi_V}{\delta \mu}(\mu) = V$$
 and  $\frac{\delta \Psi_W}{\delta \mu}(\mu) = 2W \star \mu$ .

 $\hookrightarrow$  For  $W(x) = x^2$ , we obtain  $\Psi_W = \text{Var}$ .

### Gibbs principle on path space

A concrete example

Gibbs principle

Large deviations

Statistical physic

Path space

Stochasti control

References

### Theorem (Gibbs path measure)

For any minimiser  $\overline{\mu}_{[0,T]}$  of

$$\inf_{\substack{\mu \in \mathcal{P}(\mathcal{C}([0,T],\mathbb{R}^d))\\ \forall t \in [0,T], \ \Psi(\mu_t) \leq 0}} H(\mu_{[0,T]}|\nu_{[0,T]}).$$

some  $\overline{\lambda}$  in  $\mathcal{M}_+([0,T])$  exists s.t.

$$\boxed{\frac{\mathrm{d}\overline{\mu}_{[0,T]}}{\mathrm{d}\nu_{[0,T]}}\big(x_{[0,T]}\big) = \big(Z_T^{\Psi}\big)^{-1}\exp\left[-\int_0^T \frac{\delta\Psi}{\delta\mu}(\overline{\mu}_t,x_t)\overline{\lambda}(\mathrm{d}t)\right],}$$

with  $\Psi(\overline{\mu}_t) = 0$   $\overline{\lambda}$ -a.e. Sufficient condition in the convex case.

### From conditioning to control

A concrete example

Gibbs principle

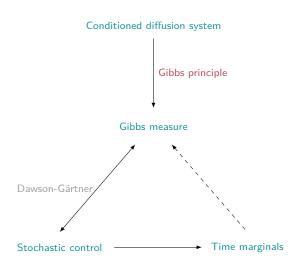
#### Stochasti control

Relation with

Time margina

Interacting

References



# Stochastic control in the diffusion setting

A concrete example

Gibbs principle

# Stochasti control

Relation with path space

Time marginals

Interacting

. .

### Controlled process

$$\begin{cases} \mathrm{d} X_s^{t,\mu,\alpha} = -\nabla V(X_s^{t,\mu,\alpha}) \mathrm{d} s + \alpha_s \mathrm{d} s + \mathrm{d} B_s, & t \leq s \leq T, \\ X_t^{t,\mu,\alpha} = X_t^{t,\mu}, & X_t^{t,\mu} \sim \mu, \end{cases}$$

for some adapted process  $\alpha = (\alpha_s)_{t \leq s \leq T}$ .

#### Value function

$$V(t,\mu) := \inf_{\substack{(lpha_s)_{t \leq s \leq T} \ orall s \in [t,T], \ \Psi(\mathrm{Law}(X_s^{t,\mu,lpha})) \leq 0}} \mathbb{E} \int_t^T rac{1}{2} |lpha_s|^2 \mathrm{d}s.$$

- $\hookrightarrow$  One looks for  $\overline{\alpha}$  which realises  $V(0,\overline{\mu}_0)$ .
- ⇒ Control problem with law constraints [Dau21].

### Pathwise law

A concrete example

Gibbs principle

Stochastic control

Relation with path space

Time marginal

Interacting particles

References

#### Theorem

There exists an optimal  $\overline{\alpha}_t = \varphi(t, X_t^{\overline{\alpha}})$  such that

$$\mathrm{Law}(X^{\overline{\alpha}}_{[0,T]}) = \overline{\mu}_{[0,T]},$$

and

$$\inf_{\substack{\mu \in \mathcal{P}(C([0,T],\mathbb{R}^d))\\ \forall t \in [0,T], \, \Psi(\mu_t) \leq 0}} H(\mu_{[0,T]}|\nu_{[0,T]}) = H(\overline{\mu}_0|\nu_0) + V(0,\overline{\mu}_0).$$

⇒ This characterises time marginals.

# Time marginals

A concrete example

Gibbs principle

Stochasti control

Relation with path space

Time margina

Interacting

Reference

### Potential mean-field game structure

 $(arphi,(\overline{\mu}_t)_t)$  is a solution of the MFG system

$$\begin{cases} \partial_t \varphi - \nabla V \cdot \nabla \varphi + \frac{1}{2} \Delta \varphi - \frac{1}{2} |\nabla \varphi|^2 = -\overline{\lambda} \frac{\delta \Psi}{\delta \mu} (\overline{\mu}_t), \\ \partial_t \overline{\mu}_t - \operatorname{div} (\overline{\mu}_t \nabla V + \overline{\mu}_t \ \nabla \varphi(t, \cdot) + \frac{1}{2} \nabla \overline{\mu}_t) = 0, \\ \varphi(T, \cdot) = 0. \end{cases}$$

As  $N \to +\infty$ , this describes a Nash equilibrium for the game

$$\inf_{\alpha^1,..,\alpha^N} \sum_{i=1}^N \mathbb{E} \int_0^T \frac{1}{2} |\alpha_t^i|^2 \mathrm{d}t + \int_0^T \frac{\delta \Psi}{\delta \mu} (\pi(\vec{X}_t^N), X_t^i) \overline{\lambda}(\mathrm{d}t).$$

### Summary

A concrete example

Gibbs principle

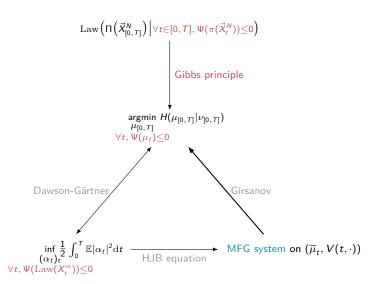
Stochasti

Relation with path space

Time marginal

Interacting particles

References



# The case of interacting particles

A concrete example

Gibbs principle

# Stochastic control

Relation with

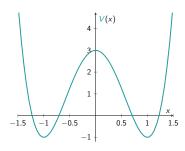
Time marginal

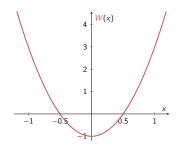
Interacting particles

References

#### Mean-field interaction:

$$\mathrm{d}X_t^{i,N} = -\nabla V(X_t^{i,N})\mathrm{d}t - \frac{1}{N}\sum_{j=1}^N \nabla W(X_t^{i,N} - X_t^{j,N})\mathrm{d}t + \mathrm{d}B_t^{i,N}$$





Confinement potential V

Interaction potential W

#### Conditioning:

$$\forall t \in [0, T], \quad \Psi(\pi(\vec{X}_t^N)) \leq 0.$$

### Mean-field stochastic control

A concrete example

Gibbs principle

# Stochasti control

Relation with path space

Time margi

Interacting

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### Controlled process

$$\begin{cases} \mathrm{d}X_{s}^{t,\mu,\alpha} = -\nabla V(X_{s}^{t,\mu,\alpha})\mathrm{d}s - \nabla (W\star \mathrm{Law}(X_{s}^{t,\mu,\alpha}))(X_{s}^{t,\mu,\alpha})\mathrm{d}s \\ + \alpha_{s}\mathrm{d}s + \mathrm{d}B_{s}, \quad t\leq s\leq T, \\ X_{t}^{t,\mu,\alpha} = X_{t}^{t,\mu}, \quad X_{t}^{t,\mu} \sim \mu, \end{cases}$$

for some adapted process  $\alpha = (\alpha_s)_{t \le s \le T}$ .

#### Value function

$$\mathcal{V}(t,\mu) := \inf_{\substack{(\alpha_s)_{t \leq s \leq T} \\ \forall s \in [t,T], \ \Psi(\mathrm{Law}(X^{t,\mu,\alpha}_s)) \leq 0}} \mathbb{E} \int_t^T \frac{1}{2} |\alpha_s|^2 \mathrm{d}s.$$

⇒ McKean-Vlasov control problem with law constraints.

### Pathwise law

A concrete example

Gibbs principle

Stochastic control

Relation with path space

Time marginal

Interacting particles

References

#### **Theorem**

There exists an optimal control such that  $\mathrm{Law}(X^{\overline{\alpha}}_{[0,T]})=\overline{\mu}_{[0,T]}$  and

$$\inf_{\substack{\mu \in \mathcal{P}(C([0,T],\mathbb{R}^d)) \\ \forall t \in [0,T], \, \Psi(\mu_t) \leq 0}} H(\mu_{[0,T]} | \Gamma(\mu_{[0,T]})) = H(\overline{\mu}_0 | \nu_0) + \mathcal{V}(0,\overline{\mu}_0).$$

### Time-marginals

$$\begin{cases} \partial_t \varphi - \nabla V \cdot \nabla \varphi - 2 \nabla (W \star \overline{\mu}_t) \cdot \nabla \varphi + \frac{1}{2} \Delta \varphi \\ - \frac{1}{2} |\nabla \varphi|^2 = -\overline{\lambda} \frac{\delta \Psi}{\delta \mu} (\overline{\mu}_t), \\ \partial_t \overline{\mu}_t - \operatorname{div} (\overline{\mu}_t \nabla V + \nabla W \star \overline{\mu}_t + \overline{\mu}_t \nabla \varphi (t, \cdot) + \frac{1}{2} \nabla \overline{\mu}_t) = 0, \\ \varphi (T, \cdot) = 0. \end{cases}$$

# Thank you!

A concrete example

Gibbs principle

# Stochasti

Relation with

Time marginals

Interacting

Reference



### References

A concrete example

Gibbs principle

Stochastic control

References

conditioning to constraints

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