

# Grushin problem method

Francis Nier, LAGA, Univ. Paris XIII

Cermics March 28th 2024

## Warning:

- Semiclassical is more than semiclassical.
- A Grushin problem is not a Grushin problem.

A reference: J. Sjöstrand and M. Zworski “Elementary linear algebra for advanced spectral problems” (Ann. Inst. Fourier (2007)).

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Hörmander FIO I (1971): *"The purpose of the present paper is not to extend the more or less formal methods used in geometrical optics but to extract from them a precise operator theory which can be applied to the theory of partial differential equations"*
- A Grushin problem is not a Grushin problem.  
The name and the notations were actually introduced in the early works of J. Sjöstrand (phD, 1973)

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## Schur comple- ment and Grushin problem

## Grushin problem and mul- tiscale analysis

## Comparison of Langevin and over- damped Langevin

- Schur complement and Grushin problem. First applications.
- Multiple wells, resonances
- Comparison of Langevin and overdamped Langevin.

The Grushin problem method can be viewed as a variation of the Schur complement formula of Linear Algebra (Numerical Analysis), like the Feshbach formula (Math. Phys.) or the Lyapunov-Schmidt method (Dynamical systems and non linear analysis).

Schur complement formula

$$\begin{pmatrix} A & B \\ C & D \end{pmatrix}^{-1} = \begin{pmatrix} 1 & -A^{-1}B \\ 0 & 1 \end{pmatrix} \begin{pmatrix} A^{-1} & 0 \\ 0 & (D - CA^{-1}B)^{-1} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -CA^{-1} & 1 \end{pmatrix},$$

whenever the block  $A$  is invertible as well as its Schur complement  $(D - CA^{-1}B)$ .

Grushin problem works in a **different way**

The **block matrix** is constructed in such a way that it is invertible and its relates the invertibility of  $A$  to the invertibility of a block of  $D$ -size (finite dimension for Fredholm theory).

Let  $\theta \mapsto A_\theta$  be a continuous map from the metric space  $(X, d)$  to  $\mathcal{L}(E; F)$ .

**Important example:**

$F$  Banach space;  $E = D(A) \subset F$  the domain of a closed operator  $A$ ;

$(X, d)$  a Banach space of relatively bounded perturbations of  $A$  with bound less than 1 ( $\Rightarrow A_\theta$  with  $D(A_\theta) = D(A) = E$ ).

### Definition

The operator  $A \in \mathcal{L}(E, F)$  is Fredholm if  $\dim \ker A = a_+ < \infty$ , if its range  $\text{Ran} A$  is closed and if  $\text{codim}(\text{Ran} A) = \dim \text{coker} A = a_- < \infty$ .

$\text{Fred}(E; F)$  set of Fredholm operators.

$\text{Ind}(A) = a_+ - a_-$ .

Remember (checked below)

- $\text{Fred}(E; F)$  is an open set of  $\mathcal{L}(E; F)$ ;
- $\text{Ind}$  is constant on connected components of  $\text{Fred}(E; F)$

Let us assume  $A_{\theta_0} \in \text{Fred}(E; F)$  (here  $(X, d) = \mathcal{L}(E; F)$ ) with  $\dim \ker A_{\theta_0} = a_+$  and  $\dim \text{coker} A_{\theta_0} = a_-$ .

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# Grushin pb and Fredholm thy

$R_- : \mathbb{C}^{a_-} \rightarrow F$  et  $R_+ : E \rightarrow \mathbb{C}^{a_+}$ ,  $R_+ \in \mathcal{L}(E; \mathbb{C}^{a_+})$ ,  $R_+|_{\ker A_{\theta_0}}$  bij.

$R_-$  (resp.  $R_+^{-1}$ ) parametrizes  $\text{coker}(A_{\theta_0})$  (resp.  $\ker(A_{\theta_0})$ ) and

$$\mathcal{A}_\theta = \begin{pmatrix} A_\theta & R_- \\ R_+ & 0 \end{pmatrix} : \begin{matrix} E \\ \oplus \\ \mathbb{C}^{a_-} \end{matrix} \rightarrow \begin{matrix} F \\ \oplus \\ \mathbb{C}^{a_+} \end{matrix}.$$

is invertible for  $\theta \sim \theta_0$  and we set

$$\mathcal{A}(\theta)^{-1} = \begin{pmatrix} E(\theta) & E_+(\theta) \\ E_-(\theta) & E_{-+}(\theta) \end{pmatrix} : \begin{matrix} F \\ \oplus \\ \mathbb{C}^{a_+} \end{matrix} \rightarrow \begin{matrix} E \\ \oplus \\ \mathbb{C}^{a_-} \end{matrix}.$$

For  $\theta = \theta_0$  solving  $\mathcal{A}_{\theta_0} \begin{pmatrix} u \\ v \end{pmatrix} = \begin{pmatrix} v \\ v_+ \end{pmatrix}$  gives

$$\begin{pmatrix} E(\theta_0) & E_+(\theta_0) \\ E_-(\theta_0) & E_{-+}(\theta_0) \end{pmatrix} = \begin{pmatrix} (A_{\theta_0}|_{\ker R_+})^{-1} \pi_{\text{Ran} A_{\theta_0}} & (R_+|_{\ker A_{\theta_0}})^{-1} \\ (R_-)^{-1} \pi_{\text{Ran} R_-} & 0 \end{pmatrix}.$$

In particular  $E_-(\theta_0)$  (resp.  $E_+(\theta_0)$ ) has the maximal rank  $a_-$  (resp.  $a_+$ ), which is an open condition ( $\Rightarrow$  the same holds true for  $E_-(\theta)$  and  $E_+(\theta)$  when  $\theta \sim \theta_0$ ).

Now by setting  $R_+ u = v_+$ , the equivalences

$$(A_\theta u = v) \Leftrightarrow \begin{cases} A_\theta u + R_- 0 & = & v \\ R_+ u & = & v_+ \end{cases} \Leftrightarrow \begin{cases} u = E(\theta)v + E_+(\theta)v_+ \\ 0 = E_-(\theta)v + E_{-+}(\theta)v_+ \end{cases},$$

lead to  $(v \in \text{Ran} A_\theta) \Leftrightarrow (E_-(\theta)v \in \text{Ran} E_{-+}(\theta))$

and  $(v = 0) \Leftrightarrow (u \in E_+(\theta) \ker E_{-+}(\theta))$ .



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$$\text{lead to} \quad (v \in \text{Ran} A_\theta) \Leftrightarrow (E_-(\theta)v \in \text{Ran} E_{-+}(\theta))$$

$$\text{and } (v = 0) \quad (u \in \ker A_\theta) \Leftrightarrow (u \in E_+(\theta) \ker E_{-+}(\theta)).$$

# Grushin problem and Fredholm theory

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$$\mathcal{A}(\theta)^{-1} = \begin{pmatrix} E(\theta) & E_+(\theta) \\ E_-(\theta) & E_{-+}(\theta) \end{pmatrix} : \begin{matrix} F \\ \oplus \\ \mathbb{C}^{a_+} \end{matrix} \rightarrow \begin{matrix} \oplus \\ \mathbb{C}^{a_-} \end{matrix},$$

$$(v \in \text{Ran} A_\theta) \Leftrightarrow (E_-(\theta)v \in \text{Ran} E_{-+}(\theta))$$

$$(u \in \ker A_\theta) \Leftrightarrow (u \in E_+(\theta) \ker E_{-+}(\theta)),$$

$$\text{rank} E_-(\theta) = a_- \quad , \quad \text{rank} E_+(\theta) = a_+ \quad \text{when } \theta \sim \theta_+.$$

$\dim \text{Ran} E_{-+}(\theta) < \infty$  and  $E_-(\theta) \in \mathcal{L}(F; \mathbb{C}^{a_-}) \Rightarrow \text{Ran} A_\theta$  closed.  
Because  $E_-(\theta)$  is onto and  $E_+(\theta)$  is one to one, we deduce

- $\dim \text{coker} A_\theta = \dim \text{coker} E_{-+}(\theta)$ ;
- $\dim \ker A_\theta = \dim \ker E_{-+}(\theta)$ .

## Proposition

With the above notations and with  $\theta \in \mathcal{V}_{\theta_0}$  small neighborhood of  $\theta_0$ ,  $A_\theta \in \text{Fred}(E; F)$  and

$$\text{Ind} A_\theta = \text{Ind} E_{-+}(\theta) = a_+ - a_-$$

which does not depend on  $\theta \in \mathcal{V}_{\theta_0}$ . Additionally when  $\text{Ind}(A_{\theta_0}) = 0$ ,  $a_+ = a_-$ , we have

$$(A_\theta \text{ invertible}) \Leftrightarrow (E_{-+}(\theta) \text{ invertible}) \Leftrightarrow (\det E_{-+}(\theta) \neq 0),$$

with the formula

$$A_\theta^{-1} = E(\theta) - E_+(\theta)E_{-+}(\theta)^{-1}E_+(\theta).$$

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# Examples of non s.a. operators

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When  $A : D(A) \rightarrow F$  is closed and the imbedding  $D(A) \rightarrow F$  is compact then either  $\sigma(A) = \mathbb{C}$  or  $\sigma(A) = \sigma_{disc}(A)$ .

For  $a = \frac{\partial_x + x}{\sqrt{2}}$  with  $D(a) = D(\mathcal{O}^{1/2})$ ,  $\mathcal{O} = \frac{-\Delta_x + x^2}{2}$ , compactly included in  $L^2(\mathbb{R}, dx)$ ,  $\sigma(a) = \sigma(a^*) = \mathbb{C}$ .

$\sigma(-\Delta_x + ix) = \emptyset$  for the complex Airy operator  $-\Delta_x + ix$  with  $D(A) = \{u \in L^2(\mathbb{R}, dx), -\Delta_x u \text{ and } xu \in L^2(\mathbb{R}, dx)\}$ .  
Subelliptic consequence:

$$\forall u \in C_K^\infty(\mathbb{R}^2), \quad \|(-\Delta_{x_1} + x_1 \partial_{x_2})u\|_{H^s} \geq C_K \|u\|_{H^{s+2/3}}$$

Weyl's theorem when  $B : D(A) \rightarrow F$  is compact the set  $\{(t, z) \in \mathbb{C} \times (\mathbb{C} \setminus \sigma_{ess}(A)), z \in \sigma(A + tB)\}$  is an analytic set. (locally  $\det(E_{-+}(z, t)) = 0$ )

When  $F = \ell^2(\mathbb{Z})$  consider the operators  $A$  et  $C \in \mathcal{L}(F)$  given by

$$(A\varphi)_n = \varphi_{n+1} \quad \text{et} \quad (C\varphi)_n = \delta_{n,0}\varphi_1$$

Then  $\sigma_{ess}(A) = \mathbb{S}^1$  and

- $\sigma_{ess}(A - tC) = \mathbb{S}^1$  for  $t \neq 1$ ;
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Weyl's theorem when  $B : D(A) \rightarrow F$  is compact the set  $\{(t, z) \in \mathbb{C} \times (\mathbb{C} \setminus \sigma_{ess}(A)), z \in \sigma(A + tB)\}$  is an analytic set. (locally  $\det(E_{-+}(z, t)) = 0$ )

When  $F = \ell^2(\mathbb{Z})$  consider the operators  $A$  et  $C \in \mathcal{L}(F)$  given by

$$(A\varphi)_n = \varphi_{n+1} \quad \text{et} \quad (C\varphi)_n = \delta_{n,0}\varphi_1$$

Then  $\sigma_{ess}(A) = \mathbb{S}^1$  and

- $\sigma_{ess}(A - tC) = \mathbb{S}^1$  for  $t \neq 1$ ;
- $\sigma(A - C) = \overline{D(0, 1)}$ .



# Examples of non s.a. operators

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When  $A : D(A) \rightarrow F$  is closed and the imbedding  $D(A) \rightarrow F$  is compact then either  $\sigma(A) = \mathbb{C}$  or  $\sigma(A) = \sigma_{disc}(A)$ .

For  $a = \frac{\partial_x + x}{\sqrt{2}}$  with  $D(a) = D(\mathcal{O}^{1/2})$ ,  $\mathcal{O} = \frac{-\Delta_x + x^2}{2}$ , compactly included in  $L^2(\mathbb{R}, dx)$ ,  $\sigma(a) = \sigma(a^*) = \mathbb{C}$ .

$\sigma(-\Delta_x + ix) = \emptyset$  for the complex Airy operator  $-\Delta_x + ix$  with  $D(A) = \{u \in L^2(\mathbb{R}, dx), -\Delta_x u \text{ and } xu \in L^2(\mathbb{R}, dx)\}$ .  
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# Quantum/semiclassical separation

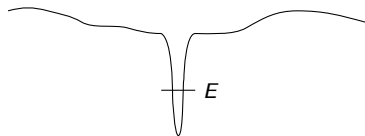
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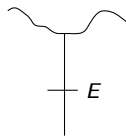
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Spectrum around  $E$  of  $-\Delta + U(x - \frac{x_0}{h}) + \tilde{V}(hx)$



or  $-h^2\Delta + U(\frac{x}{h}) + \tilde{V}(x_0 + x)$

# Quantum/semiclassical separation

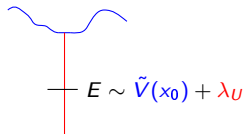
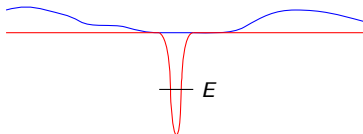
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There is an eigenvalue  $E^h = \tilde{V}(x_0) + \lambda_U + o(h^0)$  where  $\lambda_U \in \sigma(-\Delta + U)$   
The spectral elements (eigenvectors, spectral projectors. . .) of  
 $H^h = -\Delta + U(x) + \tilde{V}(hx)$  are well approximated by the ones of  $H_U$

# Quantum/semiclassical separation

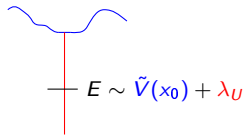
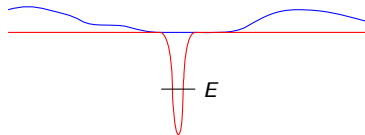
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## Exponential decay:

Quantum Hamiltonian

$$H_U = -\Delta + U(x)$$

$$H_U \psi_U = \lambda_U \psi_U, \quad \lambda_U \leq E/2 < 0$$

$$\psi_U(x) = \tilde{O}(e^{-\alpha E|x|})$$

Filled well Hamiltonian (semiclassical)

$$\tilde{H}^h = -\Delta + \tilde{V}(hx)$$

$$\tilde{V} \geq \frac{E}{4} \text{ and } \operatorname{Re} z \leq \frac{E}{2},$$

$$(z - \tilde{H}^h)^{-1}(x, y) = \tilde{O}(e^{-\alpha E|x-y|})$$

# Quantum/semiclassical separation

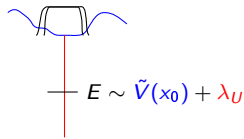
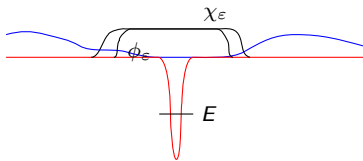
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Cut-off:  $\phi, \chi \in \mathcal{C}_0^\infty$ ,  $\phi \preceq \chi$ ,

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$$(z - \tilde{H}^h)^{-1}(x, y) = \tilde{O}(e^{-\alpha E|x-y|})$$

$$1 = \underbrace{\phi\left(\frac{hx - x_0}{\varepsilon}\right)}_{\phi_\varepsilon(hx)} + (1 - \phi)\left(\frac{hx - x_0}{\varepsilon}\right) = \underbrace{\chi\left(\frac{hx - x_0}{\varepsilon}\right)}_{\chi_\varepsilon(hx)} + (1 - \chi)\left(\frac{hx - x_0}{\varepsilon}\right)$$

## Grushin problem for $H_U$ :

$$\mathcal{A}^0(z) = \begin{pmatrix} H_U - z & R_-^0 \\ R_+^0 & 0 \end{pmatrix} : \begin{matrix} D(H_U) \\ \mathbb{C} \end{matrix} \rightarrow \begin{matrix} L^2 \\ \mathbb{C} \end{matrix}$$

With  $R_-^0 = |\psi_U\rangle$  and  $R_+^0 = \langle\psi_U|$ ,  $\mathcal{A}^0(\lambda_U)$  is invertible  $\Rightarrow$  so  $\mathcal{A}^0(z)$  is invertible for  $z \in \mathcal{V}(\lambda_U)$ :

$$\mathcal{G}^0(z) = [\mathcal{A}^0(z)]^{-1} = \begin{pmatrix} E_-^0(z) & E_+^0 \\ E_-^0 & E_{-+}^0(z) \end{pmatrix}.$$

We consider now for  $z$  close to  $\tilde{V}(x_0) + \lambda_U$

$$\mathcal{A}(z) = \begin{pmatrix} H^h - z & R_- = \chi_\varepsilon R_-^0 \\ R_+ = R_+^0 & 0 \end{pmatrix}$$

and take

$$\mathcal{F}(z) = \begin{pmatrix} \chi_\varepsilon E_-^0(z - \tilde{V}(x_0))\phi_\varepsilon + (\tilde{H}^h - z)^{-1}(1 - \phi_\varepsilon) & \chi_\varepsilon E_+^0 \\ E_-^0\phi_\varepsilon & E_{-+}^0(z - \tilde{V}(x_0)) \end{pmatrix}$$

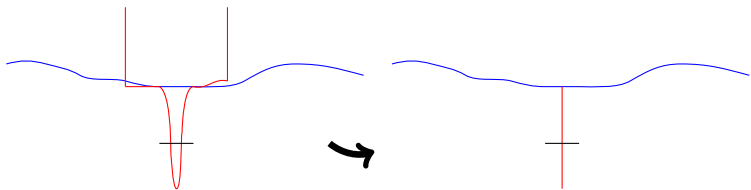
A direct computation shows (with  $h \leq \varepsilon^2$ )

$$\mathcal{A}(z)\mathcal{F}(z) = \text{Id} + \mathcal{O}(\varepsilon) + \tilde{\mathcal{O}}(e^{-\frac{\varepsilon}{h}}) \quad \text{and} \quad E_{-+}(z) = E_{-+}^0(z - \tilde{V}(x_0)) + \mathcal{O}(\varepsilon) + \tilde{\mathcal{O}}(e^{-\frac{\varepsilon}{h}})$$

Combined with  $(H^h - z)^{-1} = E(z) - E_+(z)(E_{-+}(z))^{-1}E_-(z)$  this allows to compare spectral quantities (eigenvalues, projector) around  $E_-$ .

# Changing $H_U$ , changing $\tilde{H}^h$

In the construction of the approximate inverse  $\mathcal{F}(z)$ ,  $H_U$  can be replaced by any quantum local problem (think of  $\phi_\varepsilon$  and  $\chi_\varepsilon$ -truncations)



The long range exponential decay of eigenvector with energy  $\sim E$  comes from the exponential decay estimate of  $\tilde{H}^h$ , expressed here in terms of

$$d_{Ag}(x', y'; E) = \inf_{\substack{\gamma(0)=x' \\ \gamma(1)=y'}} \int_0^1 \sqrt{|\tilde{V}(\gamma(t)) - E|} |\dot{\gamma}(t)| dt, \quad x' = hx, \quad y' = hy.$$

A perturbative analysis of  $E_{-+}(z)^{-1}$  for two different choices  $\tilde{V}_1$  and  $\tilde{V}_2$ ,  $\tilde{V}_1 = \tilde{V}_2$  in  $|h(x - x_0)| \leq R$ , leads to an error of size  $\tilde{O}(e^{-2\frac{S_R}{h}})$  where  $S_R = \inf_{|x' - x_0|=R} d_{Ag}(x', x_0)$  as long as  $\inf_{x \in \mathbb{R}} \tilde{V}_k(x) > E$ .

# Changing $H_U$ , changing $\tilde{H}^h$

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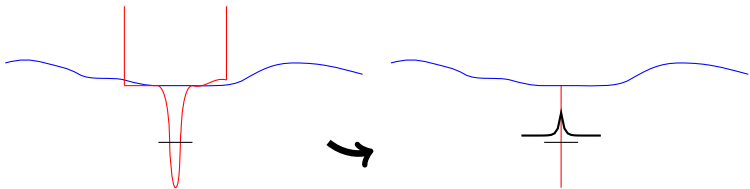
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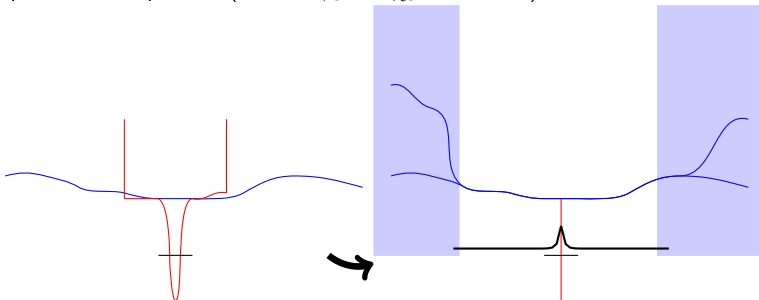
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